

Astro 507
Lecture 21
March 13, 2020

Announcements:

- Welcome to Social Distancing Cosmology!
- Lecture in real time on Zoom, video posted after
- **Preflight 4 was due: thanks for great answers**
- **Problem Set 4 due Friday after break, March 27**

Last time: began the physics of the CMB

Q: Why cosmic? Why microwave? Why background?

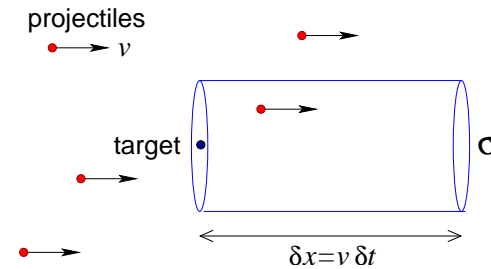
Q: where/when do CMB photons probe?

Q: physical significance of CMB sky image?

Q: when will an observer stop seeing the CMB?

Cross Section, Flux, and Collision Rate

projectiles a number density n_a
incident with speed v on targets b
with interaction cross section σ_{ab}



average collision rate per target b

$$\Gamma_{\text{per } b} = n_a \sigma_{ab} v = \sigma_{ab} j_a$$

where $j_a = n_a v$ is incident **flux**

Q: average target collision time interval?

Q: average projectile distance traveled in this time?

estimate avg time between collisions on target b :

mean free time τ

collision rate: $\Gamma = d\mathcal{N}_{\text{coll}}/dt$

so wait time until next collision set by $\delta\mathcal{N}_{\text{coll}} = \Gamma_{\text{per } b}\tau = 1$:

$$\tau = \frac{1}{\Gamma_{\text{per } b}} = \frac{1}{n_a\sigma_{ab}v} \quad (1)$$

in this time, projectile a moves distance: **mean free path**

$$\ell_{\text{mpf}} = v\tau = \frac{1}{n_a\sigma_{ab}} \quad (2)$$

no explicit v dependence, but still $\ell(E) \propto 1/\sigma_{ab}(E)$

Q: *physically, why the scalings with n, σ_{ab} ?*

Q: *what sets σ for classical billiard balls?*

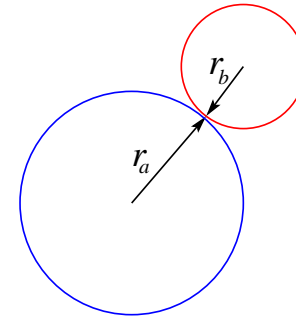
ω Q: *what set σ for $e^- + e^-$ scattering?*

Cross Section vs Particle “Size”

if particles interact only by “touching”
(e.g., classical, macroscopic billiard balls)

then $\sigma \leftrightarrow$ particle radii:

$$\sigma_{ab} = \pi(r_a + r_b)^2$$



but: if interact by force field
(e.g., gravity, EM, nuke, weak)

cross section σ *unrelated* to physical size!

For example: e^- has $r_e = 0$ (as far as we know!)

but electrons scatter via Coulomb (and weak) interaction

↳ “touch-free scattering”

Reaction Rate Per Volume

recall: collision rate *per target b* is $\Gamma_{\text{per } b} = n_a \sigma_{ab} v$
total collision rate *per unit volume* is

$$r = \frac{dn_{\text{coll}}}{dt} = \Gamma_{\text{per } b} n_b = \frac{1}{1 + \delta_{ab}} n_a n_b \sigma_{ab} v \quad (3)$$

Kronecker δ_{ab} : 0 unless particles a & b identical

Note: volume rate *symmetric* w.r.t. the two particles
as it must be

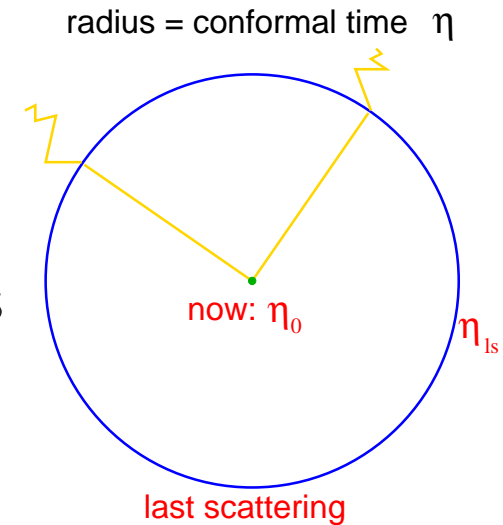
What if particles have more than one relative velocity?

CMB: Last Scattering?

CMB is a background:

all other observed sources closer

- low- z Universe transparent to CMB photons
- CMB *scattering ineffective* for these z



But scattering rate $\Gamma(\text{CMB} - \text{matter})_{\text{per}\gamma} = n_{\text{targ}}c\sigma$

- low- z U. contains atomic matter = scatterers: $n_{\text{targ}} > 0$
 - photons can and do interact with atoms/ions/electrons: $\sigma > 0$
- $\Rightarrow \Gamma(\text{CMB} - \text{matter}) > 0$: *scattering must occur!*

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Q: How can we reconcile these?

Q: Physical meaning, criterion for interaction "effectiveness"?

Particle Interactions in a FLRW Universe: Freezeouts

photon *decouple* plasma \rightarrow CMB last scattering
when: expansion redshifting & volume dilution stops interactions

$$\Gamma_{\text{scatter}} \lesssim H \quad (4)$$

or mean free time “infinite” $\rightarrow \tau \gtrsim t_H \sim t$

or mean free path “infinite” $\rightarrow \ell > d_{\text{hor,phys}}$

Q: *which of these is best to use?*

★ This criterion of very *general cosmological importance* including CMB but also all of Early Universe!

★ Since Γ depends on particle energies $\rightarrow T$ and usually Γ *increases* (strongly) with T

$\Gamma \lesssim H$ sometimes known as condition for “**freezeout**”

★ *freezeouts a central aspect of much of cosmology*

CMB, big bang nucleosynthesis, particle dark matter, 21 cm, ...

CMB Epoch: Freezeout of Cosmic Photon Scattering

Our Mission determine CMB release epoch
to do this: need photon scattering in cosmic environments

free electrons scatter photons
at low energies, cross section constant: Thomson

$$\sigma_{e\gamma} = \sigma_T = \text{const} = \frac{8\pi}{3} \left(\frac{e^2}{m_e c^2} \right)^2 = 0.665 \times 10^{-24} \text{ cm}^2$$

Q: p has same charge—why can we ignore $p - \gamma$ scattering?

Q: what is scattering rate per photon?

CMB Epoch: Egregiously Naïve Treatment

- present baryon density $n_B \approx n_e$ total electron density Q: *why?*
evolves as $n_e = n_{e,0} a^{-3}$
- using this, evaluate scattering rate *per photon*

$$\Gamma_\gamma = n_e \sigma_{TC} \stackrel{\text{naïve}}{=} n_{e,0} \sigma_{TC} a^{-3} \sim 5 \times 10^{-21} \text{ s}^{-1} a^{-3} \quad (5)$$

- expansion rate evolves roughly as matter-dom: $H = H_0 a^{-3/2}$

compare scattering and expansion rates:

$$\frac{\Gamma_\gamma}{H} \stackrel{\text{naïve}}{\simeq} 2 \times 10^{-3} a^{-3/2} = 2 \times 10^{-3} (1+z)^{3/2} \quad (6)$$

Q: *implications of $z = 0$ value?*

- this would imply $\Gamma_\gamma > H$ when $z \gtrsim 60$
Q: *what is qualitatively promising about this?*
- quantitatively, this is **wrong**: $z_{\text{last scatter}} \gg 60$
Q: *where did we go wrong?*

U. mostly composed of diffuse (gaseous) matter

Q: what are possible states of this matter?

Q: how does each interact with photons?

Q: which absorbs/scatters the most, least efficiently?

Demo: flame in projector beam

Q: brighter or darker?

Q: why do we get the result we do?

Photon Scattering Agents

Photon scatter off of charged matter: atoms, ions, electrons
mostly H (90% by number, 75% by mass)
rest is mostly He, then traces of others

molecules: H₂ essentially **invisible** Q: *why?*

neutral atoms: “H I” – essentially **invisible**

unless $E_\gamma =$ level difference, e.g., $E(\text{Ly}\alpha) = E_2 - E_1 = 10.2$ eV
or $E_\gamma > 13.6$ eV binding

ionized gas/plasma: *free e^- readily scatter photons $e\gamma \rightarrow e\gamma$*
at low energy $E_\gamma \ll m_e c^2$, Thomson scattering

$$\sigma_{e\gamma} = \sigma_T$$

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Q: *lesson for CMB*

The CMB and Recombination

In cosmic matter, photon scattering controlled by availability of **free electrons** – bound e don't count!

- ▷ ionized U: e^- abundant, scattering rapid
- ▷ neutral U: H essentially transparent to thermal background

ionized \leftrightarrow opaque

neutral \leftrightarrow transparent

CMB originates in **(re)combination**

in transition $p + e \rightarrow H + \gamma$ “the fog clears”

- plasma \rightarrow neutral H
- photon last scattering \rightarrow free streaming
- drunken stagger \rightarrow sober sprint

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Q: what (directly) determines when photons decouple from plasma?

Q: how is recombination different from decoupling? related?

Recombination and Decoupling

decoupling set by *freezeout* of scattering
as seen by photons \rightarrow when $\Gamma_{\text{scatter,per}\gamma} \lesssim H$

U. transition: **opaque \rightarrow transparent**

sets "*cosmic photosphere*" at which CMB released

(re)combination is when $p + e \rightarrow H + \gamma$

U. transition: **ionized \rightarrow neutral**

these are *logically and physical distinct* epochs

but close in time and physically *related*:

photon scattering dominated by *free e^-* : *Thomson scattering*

and free e^- abundance drops enormously at recombination

\rightarrow recombination leads to decoupling

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Q: *pre-decoupling, what should photon spectrum be?*

Q: *how are photon, plasma temperatures related?*

Cosmic Thomson Scattering

Pre-decoupled photons in thermal equilib with plasma

→ initially I_ν is Planck spectrum, $T_\gamma = T_e$

Thomson scattering continues until free e gone

Fun facts about Thomson scattering $e\gamma \rightarrow e\gamma$

- ▷ interaction strength **energy-independent**: σ_T a constant
- ▷ an **elastic** process: photon energy essentially **unchanged**
- ▷ a “two-to-two” reaction: photon number **conserved**
- ▷ scattering **anisotropic** relative to initial photon direction
angular distribution (scattering per solid angle $d\Omega$)

$$\frac{d\sigma}{d\Omega} = \frac{1}{2} \left(\frac{e^2}{m_e c^2} \right)^2 (1 + \cos^2 \theta) \quad (7)$$

includes a **quadrupole** component → creates polarization!

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Assume that recomb is a freezeout *only* of Thomson:

Q: *implications for post-recomb (i.e., observed) CMB spectrum?*

CMB Spectrum: The Magic of Thomson Scattering

Thomson implications for cosmic last scattering:

- σ_T energy-indep \rightarrow simultaneous freezeout at all freq ν
- elastic scattering \rightarrow no change in spectral *shape*
only changes photon directions
- photon number cons \rightarrow don't add or subtract to spectrum
- anisotropic scattering w.r.t. initial photon direction
but *if* initial directions isotropic \rightarrow no net anisotropy created

magic of Thomson scattering:

- ★ simultaneous freezeout of all photons (all ν)
- ★ photon spectrum *preserved*

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Q: *implication of number conservation of Thomson/Compton?*

The CMB Demands a Hot Big Bang

observe *thermal* (Planck) CMB spectrum today

⇒ *thermal* CMB spectrum *pre*-decoupling!

⇒ in early U: photons thermalized, coupled to matter!

Cosmic matter & radiation once in “good thermal contact”

→ but this requires much higher T , ρ than seen today

→ CMB demands Universe went through *hot*, *dense* early phase

⇒ **CMB** → *hot big bang*

Compton/Thomson scattering conserves photon number

but Planck spectrum has fixed number density at T

⇒ **early Universe needed photon number-changing processes**

e.g., bremsstrahlung $e + \text{nucleus} \rightarrow e + \text{nucleus} + \gamma$

moreover: we will see that $n_\gamma \sim 10^9 n_{\text{baryon}}$

⇒ need huge photon source! Q: *ideas?*

Q: *real-Universe complications?*

in the real Universe, non-Thomson processes operate

most notably: as recombination begins, *neutral H* present
resonant emission and absorption due to H lines
does lead to *non-thermal distortions* in CMB

but turns out distortions are at high frequency
i.e., nonthermal perturbations expected to be significant
only at $h\nu \gtrsim 40kT$
why this scale? we will see...

Last Scattering: Including Recombination

Recombination Revisited

For simplicity, we will assume baryons are only protons

www: laboratory hydrogen plasma

and will consider only Thomson scattering (excellent approx!)

Then: scattering rate per photon is

$$\Gamma_\gamma = n_{e,\text{free}}\sigma_T c \propto n_{e,\text{free}} \quad (8)$$

and last scattering when $\Gamma_\gamma \simeq H$

last scattering/decoupling controlled by *free electron density*

$n_{e,\text{free}}$ changes due to

- cosmic volume expansion $\propto a^{-3}$
- recombination: free e^- lost to neutral H

rewrite to account for each $n_{e,\text{free}}$ effect separately:

$$n_{e,\text{free}} = X_e n_{e,\text{tot}} = X_e n_{\text{baryon}} \quad (9)$$

- baryon density $n_{\text{b}} \propto a^{-3} \propto T^3$
gives volume dilution

- “ionization fraction”

$$X_e \equiv \frac{n_{e,\text{free}}}{n_{e,\text{free}} + n_{e,\text{bound}}} = \frac{n_p}{n_p + n_{\text{H}}} = \frac{n_p}{n_{\text{b}}} \quad (10)$$

unchanged by volume dilution

only depends on recombination thermodynamics:

i.e., $X_e = X_e(T) = X_e(z)$

in homogeneous U

Q: *what changes photon number density after recombination?*

Q: *what changes spectrum after recombination?*

Q: *naïve estimate of recombination $T_{\text{rec}}, z_{\text{rec}}$?*

Q: *zeroth-order treatment of $X_e(T)$?*

Recombination: Improved Naïve View

Given hydrogen binding energy

$$B_H = E(p) + E(e) - E(H) = 13.6 \text{ eV}$$

simple estimate of recomb epoch goes like this:

Binding sets energy scale, so

★ when particle energies above B_H : U ionized,

★ otherwise: U neutral

→ naïvely expect transition at $T_{\text{rec,naive}} = B_h \sim 150,000 \text{ K}$

But we know $T = T_0/a$, so estimate recomb at

$$\left. \begin{aligned} a_{\text{rec,naive}} &= \frac{T_0}{T_{\text{rec,naive}}} \sim 2 \times 10^{-5} \\ z_{\text{rec,naive}} &= \frac{T_{\text{rec,naive}}}{T_0} - 1 \sim 50,000 \end{aligned} \right\} \text{wrong!}$$

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Q: *guesses as to what's wrong?*

Q: *how to do this right?*