Astro ⁵⁰⁷ Lecture ²⁸ April 6, ²⁰²⁰

Announcements:

• Problem Set ⁵ due Fridaycan post questions in Homework Discussion

• Office Hours: by email, by appointment, or Instructor–after class Wednesday TA: noon-1pm Thurday (email)

Last time: BBN theory vs observation

- •• D/H very precise, ⁴He agrees measures cosmic baryons: $\eta \rightarrow n_{\textsf{B}} \rightarrow \rho_{\textsf{B}}$
- requires ² kinds of dark matter Q: how? what? $\overline{1}$
	- •• 7 Li/H strongly disagrees with D/H: "lithium problem"

Subcritical Baryons and Two Kinds of Dark Matter $0.024\leq\Omega_B\leq0.049$

 $\Omega_\mathsf{B}\ll 1$ n driv baryons do not close the universe!

 $\Omega_{\sf B} \ll \Omega_{\sf Matter} \simeq 0.3$

mne most of cosmic matter is not made of baryons!

"non-baryonic dark matter"

huge implications for particle physics–more on this to come

Measure known baryons which are directly observable optically

i.e., in *luminous* form (stars, gas): $\rho_{\text{lum}} = (M/L)_* \mathcal{L}_{\text{vis}}$ $\Omega_{\sf lum} \simeq 0.0024 h^{-1}$ $^{\texttt{1}}$ \sim 0.004 \ll $\Omega_{\texttt{B}}$

nost harvons darkl "h \Rightarrow most *baryons* dark! "**baryonic dark matter"**
O: Where are thou? Q: Where are they? \overline{M}

Where are the dark baryons?

• compact objects (white dwarfs, neutron stars, black holes) search for MACHOs: MAssive COmpact Halo Objects via gravitational microlensing www: lensing [diagram,](http://spiff.rit.edu/classes/phys240/lectures/microlens/microlens.html) MACHO event see lensing events towards LMC! but are they MACHOs or LMC stars? ...probably the latter

```
• warm/hot intergalactic medium (WHIM)<br>structure formation winfall wshock boat to
structure formation → infall → shock heat to T\sim 10^5−-10^{7}' K
note: in galaxy clusters, most baryons in
hot "intracluster" gas, not galaxies!
www: X-ray cluster
but X-rays from WHIM gas harder to see...
recent evidence of diffuse "X-ray forest"
www: Chandra spectra
```
 ω

BBN and the CMB: Battle of the Baryons

Until recently:

 $\overline{4}$

BBN was the premier means for measuring $\eta \propto \Omega_B$

 \rightarrow the best cosmic "baryometer"

Now: CMB independently measures η

battle of the baryons

compare independent measures of η test of cosmology!

```
If agreement: big bang working very well!
z∼quantitatively consistent with z \sim 10^3 theory & CMB
  \sim 10^{10} theory & light elements
```
If disagreement: ^a pressing problem!

BBN in Light of the CMB

Planck 2018: Final Data Release $\Omega_{\rm baryon, CMB}h^2=0.022298\pm 0.000214$ $\Rightarrow \eta_{\mathsf{CMB}} = (6.104 \pm 0.058) \times 10^{-10}$

- 1% precision!
- independent of BBN!

BBN vs CMB: Testing Cosmology

pillar vs pillar!

www: [Schramm](http://www.astro.uiuc.edu/classes/astr596pc/Lectures/Images/Schramm_WMAP.jpg) plot: *η*BBN ^{vs η}CMB

Concordance!

in more detail:

- 1. use η_{CMB} as input to (Std) BBN theory,
- 2. compute light elements
- 3. compare with observations

www: abundance likelihoods (BDF, Olive, Yeh, & Young 2020)

- •• D agreement perfect! ⁴He agreement excellent
- • \bullet ⁷Li tension clearer – hot research topic "lithium problem" could point to new physics!

What's up with $\frac{7}{1}$.

- observational systematics (e.g., stellar parameters)? Quite possible.(Melendez & Ramirez 2004; FOV05)
- astrophysical systematics (e.g., depletion)? but what about ⁶Li? and Li dispersion small (\lesssim 0.2 dex)...
- BBN calculation systematics: nuke reaction rates? But well measured, and can use solar neutrinos to test dominant source: 3 He(α, γ)⁷Be (CFO04)
- new physics? if so, nature kind–didn't notice till nowotherwise, would not have believed hot big bang... $\overline{}$

Particle Dark Matter

BBN and Particle Dark Matter

BBN motivates dark matter theory & searches two ways: Quantitative. $\Omega_B\ll\Omega$ m: must have non-baryonic dark matter ...and lots of it! Qualitative. BBN success at $t \sim 1$ s \rightarrow early U as physics lab
"The universe is the neer man's particle asselerater"

"The universe is the poor man's particle accelerator"

– Ya. Zel'dovich

Big implications for–and motivations from–particle physics

Q: what can we say about DM properties generally?

- Q: what can we say if DM is in particle form?lifetime, mass, interactions, quantum $\#s$?
- Q: what known particles are candidates for non-baryonic DM?Q: does particle theory offer dark matter candidates? \circ

Elementary Particle Physics and Dark Matter

Dark matter dark: no/feeble EM, strong interactions matter: behaves as nonrelativistic material $\rightarrow \rho \propto a^{-3}$, $P \ll \rho c$ 2naturally leads to hypothesis of DM as Weakly Interacting Massive Particles: WIMPs

If DM is swarms of WIMPs, what are their properties?

lifetime: must exist today $t_0 \sim 14$ Gyr \rightarrow stable or very long-lived

mass: don't know!

only know mass dens $\rho_{\mathsf{m},\mathsf{0}}$ today on cosmic, galactic scales but without also knowing $\#$ dens $n_{\text{m},0}$, can't get $m = \rho/n$ \rightarrow in fact, with specific model, from m get $n_{\rm 0}$ $\overline{0}$

Could the Dark Matter be Neutrinos?

 $interactions/quantum \#s$: BBN: dark matter not baryonic

Standard Model of particle physics does provide ^a candidate for non-baryonic DMstable $+$ massive: neutrinos; can show (PS5):

$$
\Omega_{\nu}h^2 = \frac{\sum_{\text{species}} m_{\nu}}{92 \text{ eV}} \tag{1}
$$

...but can show (β decay, ν oscillations, CMB, LSS)
— Σ species $m_\nu\lesssim$ $\lesssim 1$ eV, and so

$$
\Omega_{\nu} \sim 0.01 < \Omega_B \ll \Omega_m \tag{2}
$$

we see: *ν*s *are* non-baryonic DM but negligible contribution to density most dark matter is not neutrinos!

Dark Matter: Who Ordered That?

Dark matter isn't neutruinos. What else could it be?

no other known particle candidates are viable! i.e., DM absent from Standard Model of Particle Physics
that asseunts for all known particles an interactions that accounts for all known particles an interactions

non-baryonic DM demands physics beyond the Standard Model

particle candidates available "off the shelf" in models of physics Beyond the Standard Model i.e., particle physics models designed to explain origin of standard model features

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examples:

lightest supersymmetric particle, axion, strangelets...

Dark Matter: Cold Relics

Director's Cut below gives details. Basic idea simple:

consider a stable particle χ and $\bar{\chi}$

- in Early U: thermally created (and annihilated) as $\chi\bar{\chi}$ pairs so always $n_\chi=n_{\overline{\chi}}$
- \bullet freezes out when cold=non-relativisitic: $m_\chi \gg T_{\rm freeze} \gg T_0$
- candidate for dark matter today
- Q: if χ always in equilibrium, what is n_χ today?
- Q: but can have $n_{\chi, \mathsf{O}}=n_{\bar\chi, \mathsf{O}}>0$ today: how?
- Q: what determines n_χ today?
- $\frac{1}{\omega}$ Q: what is required for χ to be the dark matter?
-

Dark Matter: The WIMP Miracle

for $\chi\bar{\chi}$ created and annihilated in pairs

- equilibrium abundance today $(T_0 \ll m)$ is exponentially tiny /Α in equilbrium, χ annihilation never stops!
- but in expanding U: annihilations freeze out! χ particles too dilute to find each other!

Freezout condition: Γ $_{\mathsf{ann}}=n_\chi \sigma_{\mathsf{ann}} v=H$

- r ρ ρ τ $\rho \rightarrow 1$ \star smaller σ _{ann} → earlier freeze → fewer annihilations
smaller cross section → higher relic abundance smaller cross section → higher relic abundance
the weak prevaill the weak prevail!
- * today want $m_{\chi} n_{\chi,0} = \rho_{DM,0}$ implies σ_{ann} is at Weak scale: T_{freeze} ~ 1 TeV ~ E_{LHC} the WIMP miracle
- But: intenstive WIMP searches have found nothing (so far) See Director's Cut for status and outlook: **new ideas needed!** 14

Director's Cut Extras

Particle Dark Matter: Thermal RelicsKolb & Turner, Ch. 5; Dodelson Ch. 3.4

Consider stable particle species χ ($\&$ antiparticle $\bar{\chi}$)

- nonrelativistic today: $m_\chi \gg T_0 \sim 3 \times 10^{-4}$ meV
- thermally produced in the early universe

What determines its abundance today? Q : if χ is still in thermal (chemical) eq? Q: and so?

Relic Particles

for non-relativistic species: if *still in (chemical) equilibrium*: number density

$$
n_{\chi} = g_{\chi} \left(\frac{m_{\chi} T}{2\pi\hbar^2}\right)^{3/2} e^{-(m_{\chi} - \mu_{\chi})/T}
$$
(3)

chem potential: $\mu \neq 0$ iff conserved particle number

if χ number *not conserved*–i.e., equal numbers of χ and $\bar{\chi}$ then $\mu_\chi=$ 0, and so $n_\chi \sim e^{-m_\chi/T} \to 0$
→ no relic particles remain – terrible \Rightarrow no relic particles remain – terrible dark matter candidate!

Lessons: relic dark matter particles should

- either have particle/antiparticle asymmetrythis is thought to be origin of baryons
- or must have dropped out of equilibrium Q: how might this happen?

Freezeout and Relic Abundance of ^a Symmetric Species

a symmetric species χ has a cosmic abundance with *equal* numbers of particle and antiparticle \ldots or particle $=$ antiparticle

thus $n_\chi = n_{\overline{\chi}}$ exactly: no "net χ number"
A semplete applituation would leave no r \Rightarrow complete annihilation would leave no remaining particles

but: annihilation requires particle interactions! these must compete successfully with expansion & cooling

in cosmic setting, essentially *guaranteed* that at some point **annihilations freeze out**:

 $\Gamma(\chi\bar\chi\to {\rm stuff}) < H$ nonzoro rolic χ al

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 \Rightarrow nonzero relic χ abundance, mass density also guaranteed!
Or so does this quarantee that χ is the dark matter? Q: so does this guarantee that χ is the dark matter?

Annihilation Freezeout

Sketch of calculation appears here; more details in extras

Annihilation rate per χ (and $\bar{\chi}$) particle is

 $\mathsf{\Gamma}_{\mathsf{ann}}(T) \simeq n_{\chi,\mathsf{eq}}(T) \; \; \langle \sigma_{\mathsf{ann}} v \rangle \sim (m_\chi T)^3$ where σ is the annihilation cross section, $3/$ 2 $\leq e^{-}$ m_χ/T $\langle \sigma$ ann $v \rangle$ (4)and $\langle \sigma_{\textsf{ann}} v \rangle$ is a thermal average

Freezeout temperature T_{f} set by

$$
H(T_{\rm f}) \sim \frac{T_{\rm f}^2}{M_{\rm pl}} = \Gamma_{\rm ann}(T_{\rm f}) \sim (m_{\chi} T_{\rm f})^{3/2} e^{-m_{\chi}/T_{\rm f}} \langle \sigma_{\rm ann} v \rangle \qquad (5)
$$

dominated by exponential: $T_{\textsf{f}}\sim m_\chi$ so freezeout χ density is

$$
n_{\chi,\text{f}} \simeq \frac{H(T_{\text{f}} = m_{\chi})}{\langle \sigma_{\text{ann}} v \rangle} \sim \frac{m_{\chi}^2}{M_{\text{pl}} \langle \sigma_{\text{ann}} v \rangle} \tag{6}
$$

Relic Abundance and Density

relic χ abundance at freezeout $T_{\mathsf{f}}\sim m_\chi$:

$$
n_{\chi,\text{f}} \simeq \frac{H(T_{\text{f}} = m_{\chi})}{\langle \sigma_{\text{ann}} v \rangle} \sim \frac{m_{\chi}^2}{M_{\text{pl}} \langle \sigma_{\text{ann}} v \rangle} \tag{7}
$$

But we want χ abundance and mass density today note that after freeze, χ conserved! $\rightarrow n_\chi = n_{\chi,\textsf{f}}(a_\textsf{f}/a)^3$ \mathbf{V}_{α} . $\rightarrow Y_\chi\equiv n_\chi/n_\gamma$ DM/photon ra 3 $\propto T^3$ $\tilde{}\propto n_\gamma$ $_{\gamma}$ DM/photon ratio constant, set at freeze: Y_χ $\equiv \frac{n_{\chi,\text{f}}}{ }$ $n_{\gamma,\mathsf{f}}$ ∼ $\,m$ 2 $\chi^2/M_\mathsf{pl}\bra{\sigma}$ ann $v\,$ \ket{v} $m_{\mathbb{Z}}^3$ χ ∼1 $M_{\mathsf{pl}}m_\chi\left\langle\sigma$ ann $v\right\rangle$ (8)

So present number and mass densities are

$$
n_{\chi,0} = Y_{\chi} n_{\gamma_0}
$$
\n
$$
\rho_{\chi,0} = m_{\chi} n_{\chi,0} \sim \frac{1}{M_{\text{pl}} \langle \sigma_{\text{ann}} v \rangle}
$$
\n(9)

What have we shown? *if* a symmetric stable species ever created (annihilates but not decays) *then* annihilations will freeze, and inevitably have nonzero relic density today, namely

$$
\rho_{\chi,0} = m_{\chi} n_{\chi,0} \sim \frac{1}{M_{\text{pl}} \langle \sigma_{\text{ann}} v \rangle} \tag{11}
$$

This calculation is of the highest interest to particle physicists $Q: why?$

We have calculated a relic density

Q: Notable aspects about what it does, doesn't depend on?

Q: To what should it be compared?

Cold Relics: Present Abundance

 \star $\rho_{\psi,0}$ indep of m_ψ

 \star $\rho_{\psi,0}\propto 1/\sigma$: the weak prevail! Q: what sort of cross section is relevant here?

 \star To get "interesting" present density: $\Omega_\psi \sim 1 \to \! \rho_\psi \sim \rho_{\mathsf{crit}} \;$ demands a specific cross section

$$
\sigma_{\text{ann}} \sim \frac{n_{\gamma,0}}{\Omega_{\psi} M_{\text{p}\rho_{\text{crit}}}}
$$
\n
$$
\sim 10^{-38} \text{ cm}^2
$$
\n(13)

scale of the Weak interaction! $[\sigma_{\mathsf{weak}}(E\sim \mathsf{GeV})\sim 10^{-38}$ cm 2 \sum_{λ}]
]

The WIMP Miracle

Dark Matter candidate: if DM is ^a cold symmetric relic needed *annihilation cross section* is at Weak scale! corresponding energy: if $\sigma \sim \alpha/E^2$ then $\sigma \sim 10^{-36}$ cm² $= 10$ pb $\rightarrow E \sim 1$ TeV

deeper reason for DM as Weakly Interacting Massive Particle: WIMP

that weak-scale annihilations $\rightarrow \Omega_\chi \sim \Omega_{\rm nbdm}$: <mark>"WIMP Miracle"</mark>

How to find them?What if we do? What if we don't?

Director's Cut Extras

Freezeout and Relic Abundanceof ^a Symmetric Species

for conserved species ψ (chem. pot. $\mu_{\psi} \neq 0$) constant comoving number: $d(na^3) = 0$

$$
\Rightarrow \dot{n}_\psi + 3\frac{\dot{a}}{a}\,n_\psi = 0
$$

for non-conserved species: ^d(ⁿψ^a³) ⁼ qa³ dt ⁶⁼ 0, where $q = source/sink$ rate = creation/destruction rate per unit vol ⇒ $\Rightarrow \frac{\dot{n}_{\psi} + 3\frac{\dot{a}}{a} n_{\psi}}{a} = q$

assume annihilation $\psi\bar{\psi}{\rightarrow} X\bar{X}$ product X thermal, with chem. pot. $\mu_X \ll T \;\; \Rightarrow \;\; n_X = n_{\bar{X}}$

$$
q = q_{\text{net}} = q_{\text{prod}} - q_{\text{ann}}
$$
 (14)

$$
= \langle \sigma v \rangle_{\text{prod}} n_X n_{\bar{X}} - \langle \sigma_{\text{ann}} v \rangle_{\text{ann}} n_{\psi} n_{\bar{\psi}} \qquad (15)
$$

$$
= \langle \sigma v \rangle_{\text{prod}} n_X^2 - \langle \sigma v \rangle_{\text{ann}} n_\psi^2 \tag{16}
$$

 $\frac{2}{\pi}$

in equilib, Q : what condition holds for q ?

=

chem equilib: $q = 0 \Rightarrow \boxed{q_{\mathsf{prod}} = q_{\mathsf{ann}}}$ so in general

$$
\dot{n}_{\psi} + 3Hn_{\psi} = q = -\langle \sigma v \rangle_{\text{ann}} \left[n_{\psi}^2 - (n_{\psi}^{\text{eq}})^2 \right] \tag{17}
$$

and a similar expression for ψ

Change variables:

(1) use comoving coords: photon density $n_\gamma \propto T^3 \propto a^{-3}$, so put $Y = n_{\psi}/n_{\gamma}$ to remove volume dilution
then $n_{\psi} + 3\alpha/a$ $n_{\psi} = n_{\psi}$ \dot{V} then $\dot{n_{\psi}} + 3\dot{a}/a n_{\psi} = n_{\gamma} \dot{Y}$
(2) put $x = m_{\psi}/T \propto a$ (2) put $x = m_{\psi}/T \propto a$ since $t \propto 1/T^2 \propto x^2$, $dY/dt = dY/dx \dot{x} = H x dY/dx$

Then:

$$
Hx\frac{dY}{dx} = -n_{\gamma} \langle \sigma v \rangle_{\text{ann}} \left(Y^2 - Y_{\text{eq}}^2 \right) \tag{18}
$$

finally

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$$
\frac{x}{Y_{\text{eq}}} \frac{dY}{dx} = -\frac{\Gamma_A}{H} \left[\left(\frac{Y}{Y_{\text{eq}}} \right)^2 - 1 \right] \tag{20}
$$

where $\Gamma_A = n_\psi^\mathsf{eq} \left< \sigma v \right>_\mathsf{ann}$: annihil. rate

So: change in comoving ψ controlled by (1) annihil. effectiveness Γ/H (2) deviation from equil

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when \Gamma/H \gg 1<br> \Omega: what if V
    Q: what if Y < Y_{\text{eq}}? Y > Y_{\text{eq}}?
```

```
when \Gamma/H < 1Q: how does Y change?
```
 Q : how you you expect Y to evolve?

```
when Γ/H \gg 1, Y driven to Y_{\sf eq}when \Gamma/H < 1, Y change is small \rightarrow freezeout!
```

```
relic abundance at T{\rightarrow}0 or x{\rightarrow}\infty is<br>Very limit is value at freezeout
Y_\infty \simeq Y_{\text{eq}}(x_f): value at freezeout
```
Step back: How can ^a symmetric species have $n_\psi = n_{\bar\psi} \neq 0$ at $T \ll m$?

physically: expansion is key if $H = 0$, $Y_{\infty} = Y_{\text{eq}}(\infty) = 0$: \rightarrow all ψ find $\bar{\psi}$ partner, annihilate but $H \neq 0$: when U dilute enough, ψ never finds $\bar\psi$: i.e., Γ $\ll H$
ponzoro rolic abundanco nonzero relic abundance

hot relies:
$$
x_f \ll 1
$$
 $(T_f \gg m)$

\ncold relies: $x_f \gg 1$

Note: hot/cold relics refers to freezeout conditions But: hot/cold dark matter refers to structure formation criteria (namely, m vs temp $T_{\sf eq}$ at matter-rad equality)

Cold Relics: WIMPs

cold relic: non-relativistic at freezeout so $x_f = m/T_f \gg 1 \rightarrow T_f \ll m$
 \Rightarrow $m = m/T$ ($mT^2/3/2$ ⇒ \Rightarrow neq ~ $e^{-m/T} (mT)^{3/2}$
 \Rightarrow Y = $e^{-x} x^{3/2}$ ⇒ \Rightarrow Y_{eq} $\sim e^{-x}x^{3/2}$

Freezeout:

 $\Gamma_{\mathsf{ann}} = H \,\, \text{at} \,\, T = T_f \ \, \overline{G}$ \Rightarrow n eq $\langle \sigma v \rangle$ _{ann} $\sim \sqrt{G}T^2$

what needed to find value of T_f ?

To solve:

• need annihilation cross section for many models, $\langle \sigma v \rangle \propto v^n$ (S-wave: $n = 0$) $\Rightarrow \langle \sigma_{\text{ann}} v \rangle (x) = \sigma_1 c x^{n/2}$, where $\sigma_1 = \sigma (E = m)$ •• convenient rewrite $1/\sqrt{G} = M_{\text{Pl}} \simeq 10^{19} \text{ GeV}$

(Planck Mass)

set $\Gamma_{\text{ann}}(T_f) = H(T_f)$, and solve for T_f Find: $x_f \sim \ln(mM_{\text{Pl}}\sigma_1) \Rightarrow T_f = m/x_f$ So

$$
Y_{\infty} \simeq Y_{\text{eq}}(x_f) \tag{21}
$$
\n
$$
\sim \frac{x_f^{3/2}}{m M_{\text{Pl}} \sigma_1} \tag{22}
$$

 \rightarrow present relic number density

$$
n_{\psi,0} = Y_{\infty} n_{\gamma,0} = 400 \ Y_{\infty} \ \text{cm}^{-3} \tag{23}
$$

present relic mass density

$$
\rho_{\psi,0} = m n_{\psi,0} \simeq \frac{x_f^{3/2} n_{\gamma,0}}{M_{\text{Pl}} \sigma_1} \tag{24}
$$

What have we shown?

if a symmetric stable species ever created

(annihilates but not decays)

then annihilations will freeze, and

inevitably have nonzero relic density today.

This calculation is of the highest interest to particle physicists $Q: why?$

We have calculated a relic density Q: To what should this be compared?32

WIMP Searches: Accelerators

if WIMPs exist in nature ...and especially if they are supersymmetric particles likely to be found in $\sim few$ yrs at CERN Large Hadron Collider www: CERN, LHC SUSY/WIMP discovery would revolutionize particle physics and all but quarantee dark matter $=$ cold relics

Q: what would the signature be at ^a collider?What are challenges to digging it out?

Even if nature is not supersymmetric $\frac{1}{20}$ many particle theories predict new physics at ~ 1 TeV

WIMP Searches: Direct Detection

if WIMPs are DM \rightarrow dark halo full of them
local density $\delta = m n \times 0.3$ GeV cm $^{-3}$ $\frac{1}{2}$ local density $\rho=mn\sim 0.3$ GeV cm $^{-3}$ virial velocities $v_0^2\sim GM_{\rm halo}/R_{\rm halo}\sim (400\,$ km/s) 2 \Rightarrow WIMP flux $F_{\text{WIMP}} = nv_0$ \Rightarrow Look for tblueWIMP-nucleus elastic scattering – challenging!

Search using sensitive detectors: cryogenic, underground interaction: WIMP collision \rightarrow nuclear recoil
measure: effects of recoiling (E = - - 1 = 100 measure: effects of recoiling ($E_{\sf kin} \sim 1-100$ keV) nucleus Q: for example?

WIMP-nucleus recoil signatures

- ⊳ energy injection: recoil heats detector crystal specific heat $C=dE/dT\sim T^3$ $\Delta T = \Delta E/C \propto T^{-3}$ if supercold, can detect ΔT rise
- ⊳ momentum transfer: detector lattice (phonons) excited
- ⊳ scintillation, ionization: charged recoil nucleus excites medium relax via γ, e , phonon emission \rightarrow detect these

Hints at WIMPS?

several direct detection experiments see...anomalies

- DAMA ([≥] 1998): ²⁵⁰ kg NaI, Gran Sasso, Italy annual modulation seen ! very high significance Q : why is $P = 1$ yr modulation interesting?
- CRESST (2011): $CaWO₄$ crystals, 730 kg days, Gran Sasso excess of events in signal region
- CoGENT (2011, 2013): ¹⁰⁰ ^g Ge, Soudan, Minnesota annual modulation announced
- CDMS Si (2013): silicon, low-background, ¹²⁴ kg days, Soudan 36excess of events in signal region

what if anomalies are dark matter?

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www: plots of \sigma_{\chi N \to \chi N} vs m_{\chi}
```
- recoils are low-energy \rightarrow suggest "light" dark matter
 m_{max} 10 m_{max} 10 GoV; work nuclear recoil $m_{\chi}\sim 10 m_{\textsf{nucleon}}\sim 10\,$ GeV: weak nuclear recoil
- curse: low-energy recoils more difficult to dig from noise
- note: not all anomalies are consistent with each other

But: *many other experiments see nothing*, especially

- LUX: ³⁷⁰ kg liquid Xe, Sanford Laboratory, South Dakota
- SuperCDMS: SNOLab, Canada

at face value, LUX rules out other signals though alternatives remain (DM-nucleon spin dependence, DMbound states)

clearly: situation messy and confused!

that's still not all...

Q: astrophysical means infer WIMP existence and properties?

WIMP Searches: Indirect Detection

if WIMPs are DM → Galactic dark halo full of them
but Galastis halo density → sesmis mean but Galactic halo density \gg cosmic mean \rightarrow → annihilation rate $q \propto \langle \sigma_{\text{ann}} v \rangle \rho_{\text{wimp}}^2$ can be large \rightarrow annihilation products potentially observable

Local annihilations

Q: how see if $\chi \bar\chi \to \gamma \gamma$ only?
Q: how see if $\chi \bar\chi \to \gamma \tau$ other St Q: how see if $\chi \bar{\chi} \rightarrow$ other Standard Model particles?
2.8 $\chi \bar{\chi} \rightarrow e^+ e^-$ or $\chi \bar{\chi}^2$ e.g., $\chi \bar{\chi} {\rightarrow} e^+ e^-$ or $q \bar{q}$?

Galactic center annihilations

Q: how see if $\chi \bar\chi \to \gamma \gamma$ only?
Q: how see if $\chi \bar\chi \to \gamma \tau$ other St

 $\overset{\omega}{\circ}$ Q: how see if $\chi \bar{\chi} \rightarrow$ other Standard Model particles? e.g., $\chi\bar\chi{\to}e^+e^-$ or $q\bar{q}\,$?

Indirect Detection: Local Annihilation Signatures

if $\chi\bar\chi\to\gamma\gamma$ only: line emission $E_\gamma\sim m_\chi$ \to losal contribution to diffuse χ signa \Rightarrow local contribution to diffuse γ signature
but: two photon appibilation $\sqrt{2}$) seeming but: two-photon annihilation $\chi \bar{\chi} \rightarrow \gamma \gamma$ must be suppressed
also x has direct EM counling and electric charge an DM not else χ has direct EM coupling \rightarrow electric charge \rightarrow DM not dark! but *can and often do* have things like $\chi \bar{\chi} \rightarrow \pi' s \rightarrow \gamma' s$

if $\chi\bar\chi{\to}q\bar q$: hadronize, sometimes to nucleons $N\bar N$ source of $\bar{n}, \bar{p},$ and $\bar{d}=$ $|\bar{n}\bar{p}|$ \Rightarrow can look for these in cosi \Rightarrow can look for these in *cosmic rays!* but *"foreground"*: "normal" antimatter from cosmic ray propagation e.g., $p_{\textsf{cr}} + p_{\textsf{ism}} \rightarrow p p p \bar{p}$

 $4C$

if $\chi\bar\chi\!\rightarrow\!e^+e^-$: local source of *high-energy* e^+

Indirect Detection: Galactic Center Annihilation

Galactic center is ρ_{DM} peak \rightarrow annihilation goldmine!?!

Direct Photon Production

$$
★ \chi \bar{\chi} → \gamma \gamma line: E_{\gamma} = m_{\chi}, and
$$

\n
$$
★ \chi \bar{\chi} → q\bar{q} → \pi^{0} → \gamma \gamma continuum E_{\gamma} < m_{\chi}
$$

Galactic center seen in TeV range

www: [HESS](http://www.mpi-hd.mpg.de/hfm/HESS/HESS.html")

but point source too localized(?), energy spectrum ^a power-law

Galactic center in GeV range

www: Fermi sky

astrophysical foregrounds large:

- cosmic-ray interactions with ISM
- in Galactic plane $p_{\text{cr}}+p_{\text{ism}}{\rightarrow}\pi^{\text{O}}$ $\mathord{\circ} \mathord{\rightarrow} \gamma \gamma$

Daylan+ (2014): strongest claims of non-astrophysical signal centered on Galactic center, axisymmetric geometry energy spectrum $\rightarrow \chi \chi {\rightarrow} b \overline{b},\; m_{\chi} \sim$ 30 GeV

Dark Matter: Where Do We stand?

Obviously, no clear detections thus far

Current status:

accelerator and astrophysical constraints are:

competitive: both place strong constraints

on particle models for WIMPS

complementary: different methods strong in different parts

of parameter space

Upgrades coming soon on all fronts

- \rightarrow the race is on!
- \rightarrow an answer will emerge in the non-distant future! $\frac{4}{3}$

If confirmed WIMP detection:

- DM found
- need particle physics beyond Standard Model
- \star payoff big! but why asymmetrical?

modified gravity?

hidden in braneworld?