Astro <sup>507</sup> Lecture <sup>33</sup>April 17, <sup>2020</sup>

Announcements:

- Preflight <sup>6</sup> was due today
- Problem Set <sup>6</sup> due next Friday April <sup>24</sup> after this: final Problem Set due Finals Weekrecall: lowest PF and PS dropped

Last time: quantum effects in inflationToday: inflation tests and begin struture formation

 $\overline{\phantom{0}}$ 

### Inflation SpectrumStatistical Properties

 $\star$  Recall: inflaton quantum modes  $\leftrightarrow$ ↔ harmonic oscillator<br>state lkk = (a)ll2 = 0=x dominated by vacuum ↔ ground state  $\|\psi_{\text{sho}}(x)\|^2 \sim e^{-x^2/2\Delta x^2}$  $\phi_k \leftrightarrow x$  fluctuations are statistically Gaussian<br>i.e. perturbations of all sizes essure but i.e., perturbations of all sizes occur, but probability of finding perturbation of size  $\delta(R)$ on scale  $R$  is distributed as a Gaussian

 $\star$  inflaton perturbations  $\rightarrow$  reheating<br>set interior matter perturbations → radiation, matter perturbations<br>same lovels in both: "adiabatic" same levels in both: "adiabatic"

<sup>2</sup> \*\*\*\*\* All of these are bona fide predictions of inflation and are testable! Q: how?

### Inflation SpectrumSlightly Tilted Scale Invariance

recall: perturbation leaving horizon have very similar amplitudeduring inflation  $\rightarrow$  nearly same for all lengthscales  $\leftrightarrow$   $k$ <br>perturbation rms amplitude perturbation rms amplitude

$$
\delta_{\inf}^2(k) \propto k^{\alpha} \tag{1}
$$

with index  $\alpha = -6\epsilon + 2\eta \ll 1$ 

 $\omega$ 

- **★** successful inflation ⇔ slow roll ⇔  $\epsilon, \eta \ll 1$  demands<br>**nerturbation spectrum nearly independent of s** perturbation spectrum nearly independent of scalenearly "self-similar," without characteristic scale"Peebles-Harrison-Zel'dovich" spectrum
- **★ successful inflation must end**  $\rightarrow \epsilon, \eta \neq 0$ <br>demands small departures from scale-ing demands small departures from scale-invariance"tilted spectrum"

### Inflation Creates Primordial Gravity Waves

Inflaton field fluctuations are inhomogeneous perturbations to cosmic mass-energy density field

### can excite **gravitational radiation**

when fluctuations have nonzero quadrupole, *i.e.*, **tensor modes** 

- <mark>cosmic gravitational wave background</mark>
- wavelengths span all scales up to Mpc
- wave amplitude directly related to density perturbations
- waves propagate unimpeded through Universe after inflationgravity wave incident through page

effect on ring of test particles



 $\rightarrow$ 

Q: how to test?

### Searching for Primordial Gravitational Waves

- waves drive quadrupole motionintroduce CMB polarization we'll see: gravitational wave excite  $B$  modes–curl features
- In principle: direct detection possible via spacetime effects! but cosmo signal below astro events (BH, NS) not accessible to aLIGO/VIRGO, likely not LISA. www: gravitational wave signal comparison

### Testing Inflation: Status to Date

test inflation by measuring density fluctuations and their statistical properties on various scales at various epochs

CMB at large angles (large scales, decoupling)

- nearly scale invariant! woo hoo! (COBE 93)
- Gaussian distribution (COBE, WMAP, Planck) www:  $3-$ yr WMAP  $T$  [distribution](http://lambda.gsfc.nasa.gov/product/map/current/pub_papers/threeyear/parameters/images/Large/ds_f22_PPT_L.png)
- WMAP, Planck: evidence for tilt! favors large scales ("red")! Planck (2013):  $\alpha=-0.035\pm0.004$  nonzero at  $\sim9\sigma!$

These did not have to be true!

Not guaranteed to be due to inflationbut very encouraging nonetheless  $\sigma$ 

### Inflation Scorecard

As designed (postdictions) inflation solve:

- $\checkmark$  horizon problem<br> $\checkmark$  flatness problem
- $\checkmark$  flatness problem
- $\checkmark$  smoothness problem
- $\checkmark$  monopole problem

But *unexpected bonus*: structure

Thus far: observed cosmic density fields

have spectrum, statistics as predicted by inflation

- $\checkmark$  nearly scale invariant<br> $\checkmark$  gaussian statistics
- √ gaussian statistics<br>√ small tilt
- $\checkmark$  small tilt

 $\overline{\phantom{0}}$ 

Frontier: CMB polarization probes of cosmic gravity waves Stay tuned!

# Intermission: Questions?

## The Inhomogeneous Universe

Origin and Evolution of Cosmic Structure

### The Large-Scale Structure of the UniverseTheoretical and Observational Landscape

On large scales, cosmo principle an excellent approximationOn small scales, fails miserablyCosmology should explain both: now open our eyes to structur e

**Theory** Goals? tools? complications? Which scales in space, time "easy" to describe? which difficult?

#### **Observations**

 Goals? observables? complications?Which scales in space, time "easy" to measure? which difficult ?

#### Arenas for theory–observation comparison

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 Which well-matched (i.e., clear results from both)? Which poorly-matched (i.e., one or both ambiguous/difficult)?What constitutes success? When are we done?

### Large-Scale Structure: The Good, the Bad, and the Ugly

#### Structure Formation Theory

Goal: describe how small density fluctuation "seeds"grow to form structure todayTools: baryon-DM-radiation-DE particle & fluid dynamics in expanding FLRW background analytic–linearized perturb theory, idealized nonlinear models numerical–full nonlinear evolution, feedback effects Complications: nonlinear processes

(virialization, shocks, star feedback)

Degree of Difficulty:

large scales easiest–smoothest, linear perturb theory accurate

 $11$ 

smallest scales hardest-very nonlinear

#### Structure Formation Observations

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 Goal: measure growth of structures over cosmic historyTools: CMB anisotropysurveys (optical, X-ray, IR, radio,  $\gamma$ -ray...): galaxies, quasars, QSO absorption systems, lensing Complications: need for statistical completeness vs sensitivity, resolution*large scales* easy in some ways: CMB very clean galaxy, quasar statistics best over largest volumes ...but difficult in others: sensitivity, resolution lowest few independent samples of structure at largest scales "cosmic variance" (e.g., see many <sup>10</sup> Mpc regions, only one at <sup>4</sup> Gpc) reshifting, absorption present challenges only <sup>a</sup> few epochs accessiblesmall scales easy in some ways: can probe locally sample many independent regions accessible at different epochs ...but difficult in others: hard to measure at large  $z$ 

#### Comparing Theory and Observation

#### Strong Tests

well-matched at large scales: linear theory accurate, observations (esp  $\mathsf{CMB}\mathsf{)}$  clean

#### Mismatches

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Theory naturally describes density evolutiondominated by dark matter–invisible! Observations naturally look at light easiest to look at most nonlinear, baryonic systems Problem: *mass* vs *light* disconnect "bias" – rarest=largest structures easiest to see and baryons collisional, dissipative→ more spatially concentrated that DM (think halos!)<br>Iso: most light from stars—but theory of star form Also: most light from stars–but theory of star formincomplete and uncertain

 $\Rightarrow$  this is the frontier!

### Building Intuition: Spherical Collapse

consider idealized initial conditions "top hat" Universe

 $\bullet$  spherical, uniform density  $\rho$ 

workhorse

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• embedded in flat, matter-dom universewith "background" density  $\rho_{\textsf{bg}}$  ("compensated" by surroundingunderdense shell)







<sup>a</sup> nonlinear problem with analytic solution!

Given: initial density contrast  $\delta_i \ll 1$  at some  $t_i$ Want to calculate: density contrast  $\delta(t)$  lucky break–Newton's "iron sphere"/Gauss' law/Birkhoff 's: in spherical matter distribution, interior ignorant of exterior  $\Rightarrow$  overdense region evolves exactly as closed universe!

PS6: solution is parametric (cycloid)

$$
a(\theta) = \frac{a_{\text{max}}}{2} (1 - \cos \theta) \tag{2}
$$

$$
t(\theta) = \frac{t_{\max}}{\pi} (\theta - \sin \theta) \tag{3}
$$

evolution parameter: "development angle"  $\theta$ 

Q:  $a, t$  for  $\theta = 0$ ?  $\theta = \pi$ ?  $\theta = 2\pi$ ? Q: so what will this look like?



Q: describe overdensity evolution qualitatively?

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• initially expand with Universe

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- but extra gravity in overdensity slows expansion
- reach max expansion at  $t_{\text{max}}$ , then begin collapse "turnaround" epoch
- formally, collapse (to a point!) at  $t_{\text{coll}} = 2t_{\text{max}}$ •

Q: what really happens when  $t \gtrsim t_{\mathsf{coll}}$ ?

### Spherical Collapse: Fate in Real Universe

Formal spherical collpase final state: *collapse to a point!* "subuniverse" goes to big crunch!

- in reality: after turnaround, infalling matter virializes marks birth of halo as collapsed object
- Note: Brooklyn is not expanding! Nor is SS, MW, LGQ: what is criterion not to expand?

Beyond the formal solution:

- after virialized, halo still overdense
	- $\rightarrow$  neighboring shells fall in
	- $\rightarrow$  mass continues to grow by accretion!<br>in real life: mergers too
- in real life: mergers too

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# Director's Cut Extras

#### Gravity Waves: Tensor Perturbations

 $\star$  so far: only looked at density (scalar) perturbations but also tensor perturbations  $\rightarrow$  gravity waves!

what's really going on: *cosmic metric* is perturbed  $spatial$  part (in a particular coordinate system  $=$  gauge):

 $\bullet$  unperturbed  $=$  FLRW

$$
d\ell^2|_{\text{FLRW}} = a(t)^2 \left( dx^2 + dy^2 + dz^2 \right) = a(t)^2 \delta_{ij} dx_i dx_j \tag{5}
$$

with perturbations

$$
d\ell^2|_{\text{pert}} = a(t)^2 e^{2\zeta} \gamma_{ij} dx_i dx_j \tag{6}
$$

with *curvature perturbation* the *scalar* function  $\zeta(\vec{x},t)$ Q: what it its physical effect? 20

perturbed metric

$$
d\ell^2|_{\text{pert}} = a(t)^2 e^{2\zeta} \gamma_{ij} dx_i dx_j \tag{7}
$$

*curvature perturbation* scalar function  $\zeta(\vec{x},t)$  changes local volume

 $\rightarrow$  locally: isotropic stretching

tensor perturbation is, to lowest order

$$
\gamma_{ij} \approx \begin{pmatrix} 1+h_+ & h_{\times} & 0 \\ -h_{\times} & 1-h_+ & 0 \\ 0 & 0 & 1 \end{pmatrix} = \delta_{ij} + \begin{pmatrix} h_+ & h_{\times} & 0 \\ -h_{\times} & -h_+ & 0 \\ 0 & 0 & 0 \end{pmatrix}
$$
 (8)

with *two independent modes* of amplitude  $h_+, h_\times$ Q: physical effect of these modes?

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tensor perturbation is, to lowest order

$$
\gamma_{ij} \approx \delta_{ij} + \begin{pmatrix} h_{+} & h_{\times} & 0 \\ -h_{\times} & -h_{+} & 0 \\ 0 & 0 & 0 \end{pmatrix}
$$
 (9)

looks like rotation: roughly speaking preserves volumebut changes angles

moreover:  $h$  satisfies massless wave equation!  $h$   $\Leftrightarrow$  gravitational radiation<br>offect on a ring of test partic effect on <sup>a</sup> ring of test particles:



### Metric Fluctuations

tensor perturbations directly are metric perturbationwhat about the inflaton perturbations?

curvature perturbation in an invariant (coordinate independent):

$$
\zeta = \Phi + H\delta t = \Phi + H\frac{\delta\phi}{\dot{\phi}}\tag{10}
$$

 $\Phi(\vec{x}, t)$  is local gravitational potential perturbation

inflation fluctuations  $\phi$  also are metric perturbations but amplitude differs from gravity wave amplitudeby factor  $H/\dot{\phi}$ 

and thus scalar perturbation variance differs by factor

 $\sum$ 

$$
r = \frac{\Delta_h^2}{\Delta_\Phi^2} \sim \left(\frac{\dot{\phi}}{H}\right)^2 \sim \epsilon \tag{11}
$$

#### Inflationary Tensor Perturbations

variance as <sup>a</sup> function of scale (wavenumber)

$$
\Delta_h^2(k) \sim \left(\frac{V}{m_{\rm pl}^4}\right)_{aH=k} \tag{12}
$$

- evaluated at "horizon crossing"  $aH=k$
- . . . . • directly probes inflation potential  $V(\phi)$ !
- • compare to density ("scalar") perturbations: tensor-to-scalar ratio

$$
r = \frac{\Delta_h^2}{\Delta_\Phi^2} = 16\epsilon \tag{13}
$$

• for  $\epsilon \ll 1$ , expect  $r \ll 1$ : scalar dominates