Astro <sup>507</sup> Lecture <sup>5</sup>Jan 31, <sup>2020</sup>

Announcements:

 $\overline{\phantom{a}}$ 

- Happy New Year!
- Problem Set <sup>1</sup> due next Friday, Feb. <sup>7</sup>Director's Cut Extras today: magnitude scale
- Preflight <sup>1</sup> was due today–thanks!

Last time: an expanding universe

- $Q$ : how do we describe cosmic kinematics  $=$  particle motions?
- Q: what is  $a(t)$  physically? units? values?
- Q: why is <sup>a</sup> important cosmologically?
- Q: what is <sup>a</sup> "comoving" coordinate?
- Q: how should cosmic matter density  $\rho$  depend on  $a$ ?

Today: cosmic dynamics – what determines  $a(t)$ ?

## Density Evolution: Matter

definition: to cosmologist <mark>matter</mark> ≡ *non-relativistic* matter

in the non-relativistic regime:

- particle speeds  $v \ll c$ , and/or  $kT\ll mc$ <sup>2</sup> (particle rest energy)
- mass is conserved

in comoving sphere with volume  $V\propto a^3$ , mass conservation gives:

$$
dM = d(\rho V) \propto d(\rho a^3) = 0 \tag{1}
$$

gives density

$$
\mathord{\text{\rm\scriptsize{N}}}
$$

$$
\rho_{\text{non-rel}} \propto \frac{1}{V} \propto a^{-3} \tag{2}
$$

density scaling with  $a$ :

$$
\rho_{\text{non-rel}} \propto \frac{1}{V} \propto a^{-3} \tag{3}
$$

today:  $\rho_{\mathsf{matter}}(t_0) \equiv \rho_{\mathsf{m},\mathsf{0}}$ 

so at other epochs (while still non-relativistic):

$$
\rho \mathsf{m} = \rho \mathsf{m}, \mathsf{o} \, a^{-3} \tag{4}
$$

Q: what is  $\rho$ <sub>m</sub>?

### Matter Density: Time Change

matter density depends only on scale factor:

$$
\rho \mathsf{m} = \rho \mathsf{m}, \mathsf{0} \ a^{-3} \tag{5}
$$

and so

$$
\dot{\rho}_{m} = -3 \rho_{m,0} \dot{a} a^{-4} = -3H\rho_{m}
$$
 (6)

Hubble sets rate for density decrease!

Q: how must this be altered in the steady-state cosmology?

## Matter and the Steady State Cosmology

steady-state cosmology adopts perfect cosmological principle:  $\triangleright$  homogeneous  $+$  isotropic  $+$  time invariant a non-evolving universe

this demands  $\dot{\rho}=0$ : density constant but expansion carries galaxies away!  $\rightarrow$  must be new matter created to replace it<br>mass creation rate per unit volume: a mass creation rate per unit volume: q:

$$
\frac{d(\rho a^3)}{dt} = q a^3 \tag{7}
$$

$$
\dot{\rho} + 3H \rho = q \tag{8}
$$

to maintain steady state: creation rate density must be

$$
q = 3H\rho
$$
  
\n $\approx 6 \times 10^{-47} \text{ g cm}^{-3} \text{ s}^{-1} = 10^{-6} \text{ GeV}/c^2 \text{ cm}^{-3} \text{ Gyr}^{-1}$   
\nQ: implications?

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### Alternative Derivation: Fluid Picture

in fluid picture: mass conservation  $\rightarrow$  continuity equation

$$
\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \tag{9}
$$

put  $\rho = \rho(t)$  and  $\vec{v} = H\vec{r}$ :

$$
\dot{\rho} + H\rho \nabla \cdot \vec{r} = \dot{\rho} + 3\frac{\dot{a}}{a} \tag{10}
$$

$$
\frac{d\rho}{\rho} = -3\frac{da}{a}\rho\tag{11}
$$

$$
\rho \propto a^{-3} \tag{12}
$$

 $\sigma$ 

## Cosmic Forces

- **on microscale:** particles scatter, collide via electromagnetic forces (also strong and weak forces) but no net electric charges or electric currents  $\rightarrow$  no EM, strong, or weak forces on cosmo scales<br>**Prossure forces:** manifectation of random velocities
- **pressure forces:** manifestation of random velocities but pressure spatially uniform → no net pressure forces!\*<br>O: why uniform? why no not P force? (recall hydrostat  $Q$ : why uniform? why no net  $P$  force? (recall hydrostat eq)
- at large scales: only force is **gravity**
- Q: what theoretical tools needed to describe this?

<sup>∗</sup>Fine print for experts:

since  $P \propto \textsf{KE}$  density, *does* contribute to net mass-energy and thus to *gravity*,<br>this is a real offect and can be important for relativistic species with  $v \approx e$ this is a real effect and can be important for relativistic species with  $v \approx c$ <br>but even in this case, no pressure forces in the usual sense.  $\sim$  ...but even in this case, no pressure *forces* in the usual sense

## Cosmodynamics Computed

cosmic dynamics is evolution of <sup>a</sup> system which is

- gravitating,
- $\bullet$  *homogeneous*, and
- isotropic

Complete, correct treatment: General Relativity $\Rightarrow$  we will sketch this starting next week

quick 'n dirty: Non-relativistic (Newtonian) cosmologypro: gives intuition, and right answer con: involves some ad hoc assumptions only justified by GR

## Ingredients of Non-Relativistic Cosmology

Inputs: for some arbitrary cosmic time  $t$ 

- motions described by  $\vec{r}(t) = a(t) \vec{r}_0$
- $\bullet$  cosmic mass density  $\rho(t)$ , spatially uniform
- $\bullet$  cosmic pressure  $P(t)$ : in general, comes with matter but for non-relativistic matter, P not important source of energy and thus mass  $(E=mc)$ so ignore: take  $P = 0$  for now (really:  $P \ll \rho c$  $^{2}$ ) and thus gravity 2)

## thus: *gravity is only force* all cosmic matter is in "free fall"

Construction:pick arbitrary point  $\vec{r}_{\mathsf{center}} = 0$ , surround by comoving sphere, radius  $r(t)$ that moves in order to always enclosesome arbitrary but fixed mass

$$
M(r) = \frac{4\pi}{3} r^3 \rho = const \qquad (13)
$$

$$
\left(\begin{array}{c}\n\cdot & \cdot \\
\cdot & \cdot \\
\hline\n\end{array}\right)^{\rho}
$$

consider <sup>a</sup> point on the sphereQ: is it accelerated?Q: what is  $\ddot{\vec{r}} = ?$ 

### Newtonian Cosmodynamics

<sup>a</sup> point on the sphere feels acceleration

$$
\ddot{\vec{r}} = \vec{g} = -\frac{GM}{r^2}\hat{r}
$$
\n(14)

with pressure  $P = 0$ 

multiply by  $\dot{\vec{r}}$  and integrate:

$$
\dot{\vec{r}} \cdot \frac{d}{dt} \dot{\vec{r}} = -GM \frac{\hat{r} \cdot d\vec{r}/dt}{r^2} \tag{15}
$$

$$
\frac{1}{2}\dot{r}^2 = \frac{GM}{r} + K = \frac{4\pi}{3}G\rho r^2 + K\tag{16}
$$

 $Q$ : physical significance of  $K$ ? of it's sign? Q: what happens when we introduce scale factor?11

## Friedmann (Energy) Equation

introduce cosmic scale factor:  $r(t) = a(t)$   $r_0$ 

"energy" eqn: Friedmann equation

$$
H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = \frac{8\pi}{3}G\rho - \frac{\kappa c^{2}}{R^{2}a^{2}}
$$
(17)

we will see: full GR gives  $K=-2r_0^2(\kappa c^2/\kappa c^2)$  $^2_0(\kappa c^2/R^2$ 2) where

- $\bullet\;\kappa=\pm1,0$ , and
- const  $R$  is lengthscale: "*curvature*" of U.

In full GR:

⊲ Friedmann eq. holds even for relativistic matter, but  $\triangleright$  where  $\rho = \sum_{{\rm species},i} \varepsilon_i/c^2$ : mass-energy density 12

## The Mighty Friedmann (Energy) Equation

fundamental equation of the Standard Cosmology:

$$
H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3}G\rho - \frac{\kappa c^2}{R^2 a^2}
$$
 (18)

Q: why is it so important?

Q: what's <sup>a</sup> variable?Q: what's <sup>a</sup> parameter?

Q:  $a(t)$  behavior if  $K = \kappa = 0$ ? if  $\kappa \neq 0$ ?

## Dissecting Friedmann

$$
H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = \frac{8\pi}{3}G\rho - \frac{\kappa c^{2}}{R^{2}a^{2}}
$$
(19)

**variables** change with time

- <sup>a</sup>: cosmic scale factor
- $\rho$ : total cosmic mass-energy density
- **parameters** constant, fixed for all time
	- $\kappa = \pm 1$  or 0: sign of "energy" (curvature) term
	- R: characteristic lengthscale, GR  $\rightarrow$  curvature scale
- Q: how does expansion of <sup>U</sup> depend on contents of U?Q: how does expansion of <sup>U</sup> not depend on contents of U?

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 $Q:$  what about acceleration– $a$ ?

## Friedmann Acceleration Equation

Newtonian analysis gives  $\ddot{a}$  for  $P{\rightarrow}0$ ่า⊲ In full GR: with  $P\neq0$ , get Friedmann acceleration eq.

$$
\frac{\ddot{a}}{a} = -\frac{4\pi}{3}G(\rho + 3P/c^2)
$$
 (20)

#### Pressure and Friedmann

- $\star$  in "energy" (a) eq.: P absent, even in full GR
- $\star$  in acceleration eq., GR → P present, same sign as  $\rho$  $2T1$ adds to "active gravitational mass"Q: why? Q: contrast with hydrostatic equilibrium?

Friedmann energy eq is "equation of motion" for scale factor

- i.e., governs evolution of  $a(t)$ .
- To solve, need to know how  $\rho$  depends on  $a$ 
	- Q: how figure this out?

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## <sup>A</sup> Matter-Only Universe

consider a universe containing *only* non-relativistic matter Friedmann:

$$
\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{\kappa c^2}{R^2} \frac{1}{a^2}
$$
\n
$$
= \frac{8\pi G}{3}\rho_0 a^{-3} - \frac{\kappa c^2}{R^2} a^{-2}
$$
\n(22)

For  $\kappa = 0$ : "Einstein-de Sitter"

$$
(\dot{a}/a)^2 = \frac{8\pi G}{3}\rho_0 a^{-3} \tag{23}
$$

evaluate today:  $H_0^2$  0 $\epsilon_0^2 = 8\pi G \rho_0/3$ 

$$
a^{1/2}da = H_0 dt \t\t(24)
$$

$$
2/3 \, a^{3/2} \ = \ H_0 \, t \tag{25}
$$

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Q: implicit assumptions in solution?

Einstein-de Sitter:

$$
t = \frac{2}{3}a^{3/2}H_0^{-1}
$$
  
\n
$$
a = \left(\frac{3}{2}H_0t\right)^{2/3} = \left(\frac{t}{t_0}\right)^{2/3}
$$
\n(27)

Now unpack the physics:

- boundary condition:  $a = 0$  at  $t = 0 \rightarrow$  "big bang"<br>•  $a \propto t^{2/3}$  O; interpretation?
- $a \propto t^{2/3}$  Q: interpretation?
- evaluate Hubble parameter

$$
H = \frac{\dot{a}}{a} = \frac{21}{3 t} \tag{28}
$$

Q: interpretation?

• present age:

$$
t_0 = \frac{2}{3} H_0^{-1} = \frac{2}{3} t_{\mathsf{H}}
$$
 (29)

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Hubble time  $t_{\mathsf{H}}$  sets scale<br>Or note that to set in whi  $Q$ : note that  $t_{\mathbf{0}} < t_{\mathsf{H}}$ : why? Other Einstein-de Sitter fun facts:

- U. half its present age at  $a = 2^{-2}$  $\frac{2}{ }$  $3 = 0.63$
- objects half present separation (and  $8\times$  more compressed) at  $t = 2^{-3/2}$  $2t_0 = 0.35t_0$
- using measured value of  $H_0$ , calculate  $t_0 = 8.9$  Gyr but know globular clusters have ages  $t_{\textsf{gc}}$   $\gtrsim$  $\gtrsim$  12 Gyr Q: huh?

# Director's Cut Extras: The Magnitude Scale

## Star Brightness: Magnitudes

star brightness (flux) measured in **magnitude** scale magnitude = "rank" : smaller *m* → **brighter**, *more* flux<br>Serry Sorry.

Magnitudes use a logarithmic scale:

• difference of <sup>5</sup> mag is factor of <sup>100</sup> in flux:

 $m_2-m_1= \frac{2.5 \log_{10}F_2/F_1}{2}$  (definition of mag scale!)

 • mag units: dimensionless! (but usually say "mag") since always a log of ratio of two dimensionful fluxes with physical units like W/m<sup>2</sup>

What is mag difference  $m_2-m_1$ :

Q: if  $F_2 = F_1$  ?

- $Q$ : what is sign of difference if  $F_2 > F_1$ ? 20
	- Q: for equidistant light bulbs,  $L_1 = 100$ Watt,  $L_2 = 50$ Watt?

#### Apparent Magnitude

a measure of star flux  $=$  (apparent) brightness

- no distance needed
- arbitrary mag zero point set for convenience: historically: use bright star Vega:  $m(\text{Vega}) \equiv 0$ then all other mags fixed by ratio to Vega flux
- ex: Sun has apparent magnitude  $m_{\odot} = -26.74$ i.e.,  $-2.5$  log $_{10}(F_{\odot}/F_{\rm Vega}) = -26.74$ so  $F_{\rm Vega} = 10^{-26.74/2.5} F_{\odot} = 2 \times 10^{-11} F_{\odot}$
- ex: Sirius has  $m_{\text{Sirius}} = -1.45 \rightarrow \text{brighter}$  than Vega<br>so:  $E_{\text{max}} = 3.8E_{\text{max}} = 8 \times 10^{-11} E$ so:  $F_{\mathsf{Sirius}} = 3.8 F_{\mathsf{Vega}} = 8 \times 10^{-11} F_{\odot}$
- $\mu_{\text{eq}}$  ex:  $m_{\text{Polaris}} = 2.02$  Q: rank Polaris, Sirius, Vega?

 $\star$  if distance to a star is known can also compute Absolute Magnitude

abs mag  $M =$  apparent mag if star placed at  $d_0 = 10$  pc

Q: what does this measure, effectively?

#### Absolute Magnitude

absolute magnitude  $M =$  apparent mag at  $d_0 = 10$  pc

places all stars at constant fixed distance

- $\rightarrow$  a stellar "police lineup"<br>by then differences in  $F$  of
- $\rightarrow$  then differences in  $F$  only due to diff in  $L$ <br>A absolute mag effectively measure lumines
- $\rightarrow$  absolute mag effectively measure luminosity

```
Sun: abs mag M_{\odot}= 4.76 mag
Sirius: M_{\mathsf{Sirius}} = +1.43 mag
Vega: M_{\rm Vega}=+0.58 mag
Polaris: M<sub>Polaris</sub> = −3.58 mag<br>- Eridani: M
\epsilon Eridani: M_{\epsilon E}ri = +6.19 mag (nearest exoplanet host; d = 3.2 pc)
Q: rank them in order of descending L?
```
Immediately see that Sun neither most nor least luminous star around23

### Distance Modulus

take ratio of actual star flux vs "lineup" fluxat abs mag distance  $d_0 = 10$  pc:

$$
\frac{F}{F_0} = \frac{L/4\pi d^2}{L/4\pi d_0^2}
$$
\n(30)

which, after simplification, leads to

$$
m - M = 5 \log \left( \frac{d}{10 \text{ pc}} \right) \tag{31}
$$

- $\bullet$  depends only on distance  $d$ , not on luminosity! can use as measure of distance
- $m M \equiv$  "distance modulus", sometimes denoted  $\mu$

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